



Addendum: A dispersive analysis of $\eta' \rightarrow \pi^+\pi^-\gamma$ and $\eta' \rightarrow \ell^+\ell^-\gamma$

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Abstract In this addendum to Ref. [1] we show that the mismatch between the ρ - ω mixing parameter $\epsilon_{\rho\omega}$ as extracted from $\eta' \rightarrow \pi^+\pi^-\gamma$ and $e^+e^- \rightarrow \pi^+\pi^-$ can be resolved by including higher orders in the expansion in e^2 in the description of the $\eta' \rightarrow \pi^+\pi^-\gamma$ decay. We repeat the analysis in this extended framework and update the numerical results accordingly.

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1 Extended formalism

Following the notation from Ref. [1] throughout, the spectrum for $P \rightarrow \pi^+\pi^-\gamma$ can be expressed as

$$\frac{d\Gamma(P \rightarrow \pi^+\pi^-\gamma)}{ds} = 16\pi\alpha\Gamma_0|F_\pi^V(s)|^2 \left| P(s)(1 + \Pi_\pi(s)) - \frac{e^2 F_{P\gamma\gamma}}{s} - \frac{g_{P\omega\gamma}}{g_{\omega\gamma}} \frac{\epsilon_{\rho\omega} - e^2 g_{\omega\gamma}^2}{M_\omega^2 - s - iM_\omega\Gamma_\omega} \right|^2, \quad (1.1)$$

generalizing Eq. (D.14) in Ref. [1] by the next order in the expansion in e^2 (the sign convention is such that $g_{P\omega\gamma} < 0$). The most important change, numerically, concerns $\epsilon_{\rho\omega} \rightarrow \epsilon_{\rho\omega} - e^2 g_{\omega\gamma}^2$ in the numerator of the ω propagator, corresponding to the photon contribution in $\epsilon_{\rho\omega}$ as defined in resonance chiral perturbation theory [2–4]. In our formalism, $\epsilon_{\rho\omega}$, determined from a fit to the bare cross section for $e^+e^- \rightarrow \pi^+\pi^-$, does not include this VP effect, in line with the definition in Ref. [5] (numerically, it evaluates to $e^2 g_{\omega\gamma}^2 = 0.34(1) \times 10^{-3}$). This shift removes the tension

observed between $\eta' \rightarrow \pi^+\pi^-\gamma$ and $e^+e^- \rightarrow \pi^+\pi^-$ in Ref. [1].

The coefficients appearing in Eq. (3.9) of Ref. [1] are generalized according to Eq. (1.1):

$$\begin{aligned} \mathcal{A}_2 &= -\Gamma(\eta' \rightarrow \pi^+\pi^-\gamma) + 16\pi\alpha \int_{4M_\pi^2}^{M_{\eta'}^2} ds \Gamma_0 |F_\pi^V(s)|^2 \\ &\quad \times \left| \frac{g_{\eta'\omega\gamma}}{g_{\omega\gamma}} \frac{\epsilon_{\rho\omega} - e^2 g_{\omega\gamma}^2}{M_\omega^2 - s - iM_\omega\Gamma_\omega} + \frac{e^2 F_{\eta'\gamma\gamma}}{s} \right|^2, \\ \mathcal{A}_1 &= 32\pi\alpha \int_{4M_\pi^2}^{M_{\eta'}^2} ds \Gamma_0 |F_\pi^V(s)|^2 \operatorname{Re} \left[P_{\text{ev}}(s)(1 + \Pi_\pi^*(s)) \right. \\ &\quad \left. \times \left(\frac{g_{\eta'\omega\gamma}}{g_{\omega\gamma}} \frac{e^2 g_{\omega\gamma}^2 - \epsilon_{\rho\omega}}{M_\omega^2 - s - iM_\omega\Gamma_\omega} - \frac{e^2 F_{\eta'\gamma\gamma}}{s} \right) \right], \\ \mathcal{A}_0 &= 16\pi\alpha \int_{4M_\pi^2}^{M_{\eta'}^2} ds \Gamma_0 |F_\pi^V(s)|^2 P_{\text{ev}}^2(s) |1 + \Pi_\pi(s)|^2. \end{aligned} \quad (1.2)$$

In the following, we provide the updated numerical results when including the additional e^2 effects as given in Eq. (1.1), implemented in the fit via Eq. (1.2).

2 Numerical results

The updated fit parameters are collected in Table 1, Fig. 1, and Table 2. The main difference to the results presented in Ref. [1] is that the shift $\epsilon_{\rho\omega} \rightarrow \epsilon_{\rho\omega} - e^2 g_{\omega\gamma}^2$ removes the tension between $e^+e^- \rightarrow \pi^+\pi^-$ and the $\eta' \rightarrow \pi^+\pi^-\gamma$ spectrum, markedly improving the quality of the combined fit.

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Table 1 Comparison of the fit outcome of the differential decay width in Eq. (1.1) to the BESIII $\eta' \rightarrow \pi^+\pi^-\gamma$ spectrum [6] of the binned maximum likelihood and minimum χ^2 strategies. The χ^2/dof is 1.30 and 1.31, respectively, with the one of the Likelihood method extracted by means of the approximation described in App. C of Ref. [7]

Quantity	Likelihood	χ^2
A [GeV^{-3}]	17.12(35)	17.09(32)
β [GeV^{-2}]	0.714(55)	0.723(45)
γ [GeV^{-4}]	-0.412(55)	-0.420(45)
$\epsilon_{\rho\omega} \times 10^3$	1.998(67)	1.997(54)
M_ω [MeV]	782.99(33)	783.00(27)

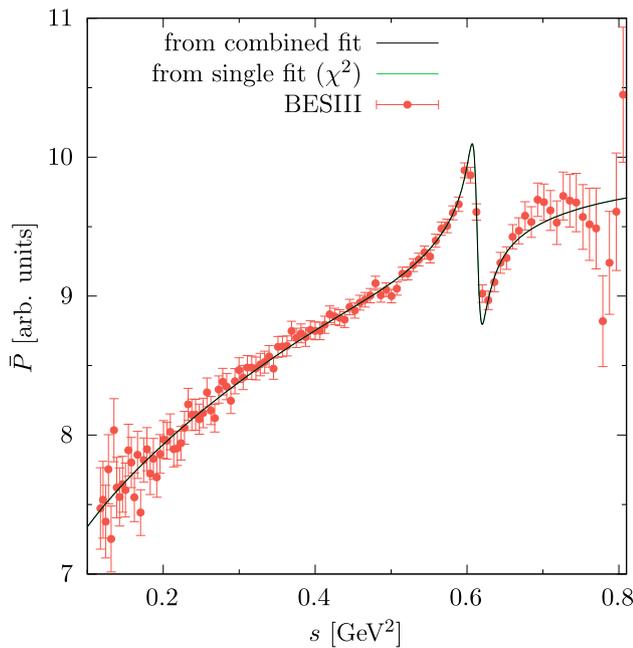


Fig. 1 Fit to the differential decay rate of $\eta' \rightarrow \pi^+\pi^-\gamma$ (individually or combined with the VFF). To highlight potential differences in the ρ - ω region, we show the associated function \bar{P} , as defined in Eq. (3.11) of Ref. [1], compared to the experimental data from BESIII [6]. The two fits cannot be distinguished on this scale

Table 2 Combined fit to several pion VFF data sets (BaBar, KLOE, CMD-2, SND) and $\eta' \rightarrow \pi^+\pi^-\gamma$ spectrum (BESIII) with overall $\chi^2/\text{dof} = 1.46$. In the row for KLOE, the three values for M_ω refer to

	χ^2/dof	M_ω [MeV]	A [GeV^{-3}]	β [GeV^{-2}]	γ [GeV^{-4}]	$\alpha_\pi \times 10^2$ [GeV^{-2}]	$\epsilon_{\rho\omega} \times 10^3$
BaBar	1.26	781.875(82)	17.10(32)	0.720(46)	-0.418(46)	5.74(14)	2.007(10)
KLOE	1.61	781.65(12)					
		782.10(17)					
		781.84(27)					
CMD-2	2.18	782.131(68)					
SND	2.16	781.457(97)					
BESIII	1.31	783.00(28)					

The updated results for the TFF are shown in Fig. 2 and Table 3. In particular, the prediction for the slope parameter

$$b_{\eta'} = 1.431(23) \text{ GeV}^{-2} \tag{2.1}$$

is reduced by about 1σ , which traces back not to the change in $\epsilon_{\rho\omega}$ (which is marginal given the fact that the fit is dominated by $e^+e^- \rightarrow \pi^+\pi^-$), but to a stronger curvature in the polynomial $P(s)$ (the coefficient γ of the quadratic term increases by a factor 3).

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the combinations of the global KLOE ω mass and the corresponding mass shifts of the three underlying data sets from 2008, 2010, 2012, respectively

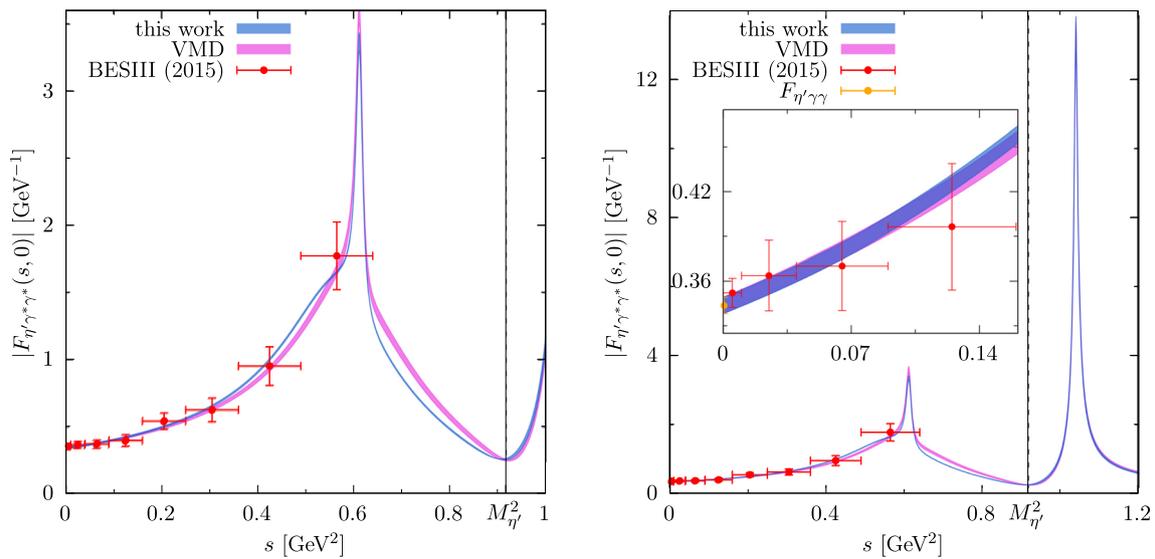


Fig. 2 Determination of the η' TFF in comparison to data from BESIII [8] (statistical and systematic errors added in quadrature) scaled with $F_{\eta'\gamma\gamma}$ and the VMD model from Ref. [1] for the ϕ resonance; for

the kinematic range accessible in η' decays (left) and a larger time-like region including the ϕ resonance with inset magnifying the low- s region (right)

Table 3 Contributions from the various components of the TFF to the sum rules of the normalization and the slope parameter

	$(I = 1)_{\epsilon_{\rho\omega}=0}$	$\Delta(I = 1)_{\epsilon_{\rho\omega}\neq 0}$	$(I = 0)_{\epsilon_{\rho\omega}=0}^\omega$	$\Delta(I = 0)_{\epsilon_{\rho\omega}\neq 0}^\omega$	$(I = 0)^\phi$	Total
Norm [%]	69.18(86)	-0.1388(19)	7.06(22)	-0.1397(47)	15.85(61)	91.9(1.1)
$b_{\eta'}$ [GeV ⁻²]	1.160(23)	0	0.1176(32)	0	0.1526(53)	1.431(23)

References

1. S. Holz, C. Hanhart, M. Hoferichter, B. Kubis, Eur. Phys. J. C **82**, 434 (2022). [arXiv:2202.05846](https://arxiv.org/abs/2202.05846) [hep-ph]
2. R. Urech, Phys. Lett. B **355**, 308 (1995). [arXiv:hep-ph/9504238](https://arxiv.org/abs/hep-ph/9504238)
3. J. Bijnens, P. Gosdzinsky, Phys. Lett. B **388**, 203 (1996). [arXiv:hep-ph/9607462](https://arxiv.org/abs/hep-ph/9607462)
4. J. Bijnens, P. Gosdzinsky, P. Talavera, Nucl. Phys. B **501**, 495 (1997). [arXiv:hep-ph/9704212](https://arxiv.org/abs/hep-ph/9704212)
5. G. Colangelo, M. Hoferichter, B. Kubis, P. Stoffer, JHEP **10**, 032 (2022). [arXiv:2208.08993](https://arxiv.org/abs/2208.08993) [hep-ph]
6. M. Ablikim et al. [BESIII Collaboration], Phys. Rev. Lett. **120**, 242003 (2018). [arXiv:1712.01525](https://arxiv.org/abs/1712.01525) [hep-ex]
7. S. Holz, Rheinische Friedrich-Wilhelms-Universität Bonn, PhD thesis (2022). <https://hdl.handle.net/20.500.11811/10336>
8. M. Ablikim et al. [BESIII Collaboration], Phys. Rev. D **92**, 012001 (2015). [arXiv:1504.06016](https://arxiv.org/abs/1504.06016) [hep-ex]